

Quantum field theory corresponding to photon driven - dissipative QPE

$$\begin{aligned}
 H = \int dz & \left[ \hat{E}^\dagger(z) \left[ \omega_0 - \frac{\hbar v^2}{2m} + \delta\omega(z) \right] \hat{E}(z) + \right. \\
 & + \frac{\partial \mu}{2} \hat{E}^\dagger(z) \hat{E}^\dagger(z) \hat{E}(z) \hat{E}(z) + \\
 & \left. + \mathbb{K} E_{nc}(z, t) \hat{E}^\dagger(z) + \mathbb{K}^* E_{nc}^\dagger(z, t) \hat{E}(z) \right] + \\
 & + \text{losses (+ amplification)}
 \end{aligned}$$

Restrict to ohmic drive

External potential  $\delta\omega(z) = \delta\omega_0 (\cos k_x x + \cos k_y y)$

lattice in both directions

↳ approx to 1 state per site

⇒ generalized Bose-Hubbard Hamiltonian

$$\begin{aligned}
 H_{BH} = \sum_{\langle ij \rangle} & -J \hat{a}_i^\dagger \hat{a}_j + \sum_i \left[ \frac{U}{2} \hat{a}_i^\dagger \hat{a}_i^\dagger \hat{a}_i \hat{a}_i + \omega_0 \hat{a}_i^\dagger \hat{a}_i + \right. \\
 & \left. + F_i(t) \hat{a}_i^\dagger + F_i^*(t) \hat{a}_i \right] + \text{losses}
 \end{aligned}$$

$J$  = tunneling (or hopping) amplitude  
 $U$  = charging energy  
 $F_i$  = driving  
 $\hat{a}_i$  = destruction operator on wf of site  $i$

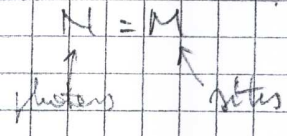


Fermionic system

$$H = \sum_i \omega_i a_i^\dagger a_i + \frac{U}{2} \sum_{\langle i, j \rangle} a_i^\dagger a_j^\dagger a_i a_j +$$

$$= J \sum_{\langle i, j \rangle} a_i^\dagger a_j$$

Restrict to consensurate occupation



\* for  $\frac{U}{J} \ll 1$ : ground state  $| \psi_g \rangle = \frac{1}{\sqrt{N!}} \left( \sum_i \frac{1}{\sqrt{N}} a_i^\dagger \right)^N | vac \rangle$

"muffled state", extended coherence in space

$$g^{(2)}(r, r') = 1$$

$$Energy = N \left[ -J \cdot c + \frac{U}{2} \right] \quad (c = \# \text{ of neighboring sites})$$

\* for  $\frac{U}{J} \gg 1$  ground state  $| \psi_g \rangle = \prod_i a_i^\dagger | vac \rangle$

"Mott-insulator" state, fixed occupation,

no coherence  $g^{(2)}(r, r') = \delta_{r, r'}$

$$Energy = 0$$

↳ Answer for  $\frac{U}{J} = \frac{U}{J_{crit}}$

Experiment with atoms

Greiner et al. Nature 415, 39 (2002)

Proposals with photons

Angelakis et al. PRA 75, 031805 (2007)

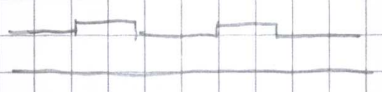
Jeanneau et al. Nat. Phys. 2, 855 (06)

Hartmann et al. Nat. Phys. 2, 869 (05)



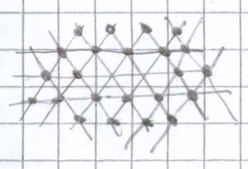
Most promising physical realizations:

a) - Planar microcavities + lateral wetting



nonlinearity  $\rightarrow$  QW exciton or  
Q. dot (ph. localization of dot  
+ inhomog. broadening)  
(analogously  $\rightarrow$  microbillars)

b) - Photonic crystal cavity:

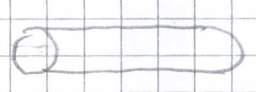


wrong holes in periodic array  $\rightarrow$  defect state  
with sub- $\lambda$  localization, high Q

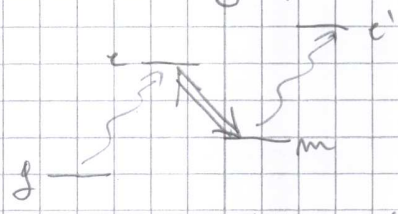
nonlinearity  $\rightarrow$  QW or Q dots  
 $\hookrightarrow$  expt. observed entanglement

$\Rightarrow$  b) - Losses for solid thin-film / evaporative such in g.s.  
 $\hookrightarrow$  dimension/density to be included

c) - hollow fibers + atoms (inside or at surface)



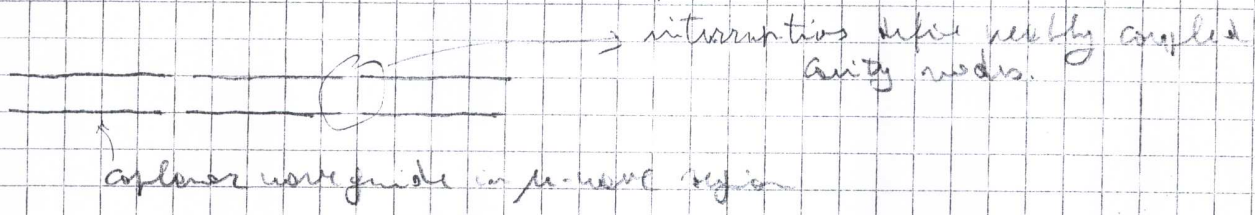
single-mode propagation  $\rightarrow$  1D  
nonlinearity from atoms, e.g. EIT



ph. not yet high enough optical thickness



d) superconducting circuits



nonlinearity: Josephson junction circuit  $\rightarrow$  effective 2-level system  
(more in Haroche's lectures)

advantages  $\frac{Q}{f} \gtrsim 100$  (while  $\sim 5$  for semiconductors)

disadvantages  $T < 1K$  (hard cryogenics)  
no single photon detector  
has to use linear amplifiers  $\rightarrow$  cannot detect excitations  
 $\rightarrow$  quantum backaction observed.

Tonks-Lieberson gas in 1D (lattice  $N \ll M$  or continuum 1D)

$\frac{U}{J} = \infty$  impenetrable photons "fermionized"

$\psi_B(x_i, \dots, x_N) = 0$  whenever  $x_i = x_j$ , fully symmetric wave exchange

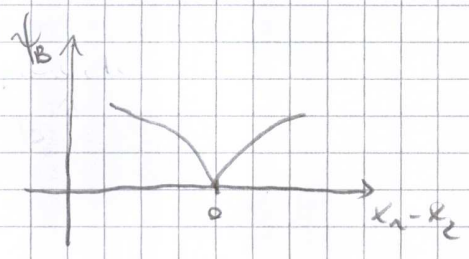
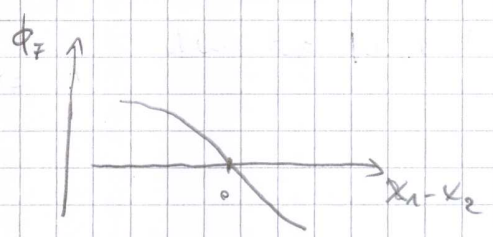
Jordan J. Math. Phys. 1, 516 ('60)

$\psi_B(x_1, \dots, x_N) = \epsilon_{x_1, \dots, x_N} \cdot \Phi_F(x_1, \dots, x_N)$  for each eigenstate

impenetrable bosons  $\rightarrow$  non-interacting fermions  
 $\rightarrow$  parity of ordering to  $x_1 < x_2 < \dots < x_N$



ground state  $\rightarrow$  Fermi sphere in  $\phi_F$



when two particles approach:

$$\left\{ \begin{array}{l} \phi_F \sim x_1 - x_2 \\ \phi_B \sim |x_1 - x_2| \end{array} \right.$$

To observe TG states in presence of disorder:

- $\rightarrow$  each eigenstate gives peak in transmission spectrum
- $\rightarrow$  structure of eigenstate deduced from coherence functions of secondary emission

PRL 103, 033601 (103)

New directions  $\rightarrow$   
(theory)

more complex many-body states  
 interplay of amplification + strong interactions  
 effect of artificial gauge field  $J \rightarrow J e^{i d_{ij}}$   
 Laughlin states, fractional statistics...

(experiment)

simulate a strongly correlated photon gas