

## Exercise 5

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1) incident state  $|\psi\rangle = \nu_a(\omega) |a; n\rangle + \nu_b(\omega) |b; n\rangle$

final state  $\hat{U}|\psi\rangle = \nu_c(\omega) |c; n\rangle + \nu_d(\omega) |d; n\rangle$

with

$$\begin{pmatrix} \nu_c(\omega) \\ \nu_d(\omega) \end{pmatrix} = e^{-i\omega(t_f - t_i)} S(\omega) \begin{pmatrix} \nu_a(\omega) \\ \nu_b(\omega) \end{pmatrix}$$

$$\langle c; n | \hat{U} | \psi \rangle = \langle c; n | \hat{U} | a; n \rangle \cdot \nu_a(\omega) +$$

$$+ \langle c; n | \hat{U} | b; n \rangle \nu_b(\omega) = \nu_c(\omega)$$

$\Rightarrow \langle (c, d); n | \hat{U} | (a, b); n \rangle$  are matrix elements of  $S(\omega)$ .

2)  $a_n^+ \dots d_n^+$  create a photon in same state.

B-S does not create/destroy photons

$\Rightarrow a_n^+, b_n^+$  go into  $c_n^+, d_n^+$  under  $U$

$$U a_n^+ U^\dagger |vac\rangle = U a_n^+ |vac\rangle \quad (\text{vacuum does not evolve})$$

$$\parallel$$
$$(r_c(\omega) c_n^+ + r_d(\omega) d_n^+) |vac\rangle$$

$$\parallel$$
$$(t(\omega) c_n^+ + r(\omega) d_n^+) |vac\rangle$$

$$\Rightarrow \begin{cases} U a_n^\dagger U^\dagger = [t(\omega) c_n^\dagger + i r(\omega) d_n^\dagger] e^{-i\omega(t_f - t_i)} \\ U b_n^\dagger U^\dagger = [t(\omega) d_n^\dagger - i r(\omega) c_n^\dagger] e^{-i\omega(t_f - t_i)} \end{cases}$$

$$3) \begin{cases} U a_n U^\dagger = [t(\omega) c_n - i r(\omega) d_n] e^{i\omega(t_f - t_i)} \\ U b_n U^\dagger = [t(\omega) d_n - i r(\omega) c_n] e^{i\omega(t_f - t_i)} \end{cases}$$

$$a_n = U^\dagger U a_n U^\dagger U = [t(\omega) U^\dagger c_n U - i r(\omega) U^\dagger d_n U] e^{i\omega(t_f - t_i)}$$

$$b_n = U^\dagger U b_n U^\dagger U = [t(\omega) U^\dagger d_n U - i r(\omega) U^\dagger c_n U] e^{i\omega(t_f - t_i)}$$

4) Inverting these relations:

$$\begin{pmatrix} U^\dagger c_n U \\ U^\dagger d_n U \end{pmatrix} = \begin{pmatrix} t(\omega) & i r(\omega) \\ i r(\omega) & t(\omega) \end{pmatrix} \begin{pmatrix} a_n \\ b_n \end{pmatrix} \cdot e^{-i\omega(t_f - t_i)}$$

5)

$$|\psi\rangle = \exp\left[-\frac{1}{2}(\alpha^2 + \beta^2)\right] \exp(\alpha a^\dagger + \beta b^\dagger) |vac\rangle$$

coherent state  $\alpha|\psi\rangle = \alpha|1\rangle, \beta|\psi\rangle = \beta|1\rangle$

$$\begin{aligned} c U |\psi\rangle &= U U^\dagger c U |\psi\rangle = U (t a + i r b) |\psi\rangle = \\ &= U \cdot (t\alpha + i r\beta) |\psi\rangle = (t\alpha + i r\beta) U |\psi\rangle \end{aligned}$$

$\Rightarrow U|\psi\rangle$  is a coherent state for c and d

Amplitude  $\gamma = t\alpha + i r\beta$  as in classical scatt

6)  $|\psi\rangle = a^\dagger b^\dagger |vac\rangle$

$$\begin{aligned} U|\psi\rangle &= U a^\dagger b^\dagger |vac\rangle = U a^\dagger U^\dagger U b^\dagger U^\dagger U |vac\rangle \\ &= (t c^\dagger + i r d^\dagger) (n a c^\dagger + t d^\dagger) |vac\rangle \\ &= [n r t c^{\dagger 2} + n r t d^{\dagger 2} + (t^2 - r^2) c^\dagger d^\dagger] |vac\rangle \end{aligned}$$

for 50/50 beam splitter  $t^2 = r^2 = 1/2$

$$\Rightarrow U|\psi\rangle = \frac{i}{2} (c^{\dagger 2} + d^{\dagger 2}) |vac\rangle =$$

$$= \frac{i}{\sqrt{2}} \left( \frac{c^{\dagger 2}}{\sqrt{2}} + \frac{d^{\dagger 2}}{\sqrt{2}} \right) |vac\rangle$$

2 particles in c

2 particles in d

7)

$$|\psi\rangle = \int \frac{dn}{2\pi} \int \frac{dn'}{2\pi} \phi(n) \phi(n') e^{-in x_a} e^{-in' x_b} a_n^\dagger b_{n'}^\dagger |vac\rangle$$

for  $\phi$  real  $\rightarrow$  wave packets centered at  $x_a, x_b$

$$U|\psi\rangle = \dots + \int \frac{dn}{2\pi} \int \frac{dn'}{2\pi} \phi(n) \phi(n') e^{-in x_a} e^{-in' x_b}$$

$$\cdot [t(n) + (n') c_n^\dagger d_{n'}^\dagger - r(n) r(n') c_{n'}^\dagger d_n^\dagger] e^{-i(n+n')(x_a-x_b)} \cdot |vac\rangle$$

Assume beam splitter is frequency independent and 50/50

$$= \dots + \frac{1}{2} \int \frac{dn}{2\pi} \int \frac{dn'}{2\pi} \left( e^{-in \bar{x}_a} e^{-in' \bar{x}_b} - e^{-in \bar{x}_b} e^{-in' \bar{x}_a} \right) \phi(n) \phi(n')$$

$$\text{with } \bar{x}_a = x_a - ct, \bar{x}_b = x_b - ct \quad \cdot c_n^\dagger d_{n'}^\dagger |vac\rangle$$

$$= \dots + \frac{1}{2} \int \frac{dk}{2\pi} \int \frac{dk'}{2\pi} \phi(k) \phi(k') \cdot e^{-i(\hbar k k')(x_0 - ct + x_0 - ct)/2}$$

$$\cdot \left( e^{-i(\hbar k k')(x_0 - x_0)/2} + e^{i(\hbar k k')(x_0 - x_0)/2} \right) \cdot \langle \dots \rangle$$

if  $|(\hbar k k')(x_0 - x_0)| \ll 1 \rightarrow$  two terms cancel each other out only cc or dd.

if  $|(\hbar k k')(x_0 - x_0)| \gtrsim 1 \rightarrow$  also col states possible.

\*  $|k - k'|$  at most of the order of  $\Delta k$ .

\*  $\Delta k \sim \frac{1}{L}$  with  $L =$  wave packet length.

$\Rightarrow \left| \frac{x_0 - x_0}{L} \right| \ll 1$  is condition for wavepacket overlap.

So: quantum interference forbids col states if incident wavepackets overlap.

Effect can not be observed with classical fields.