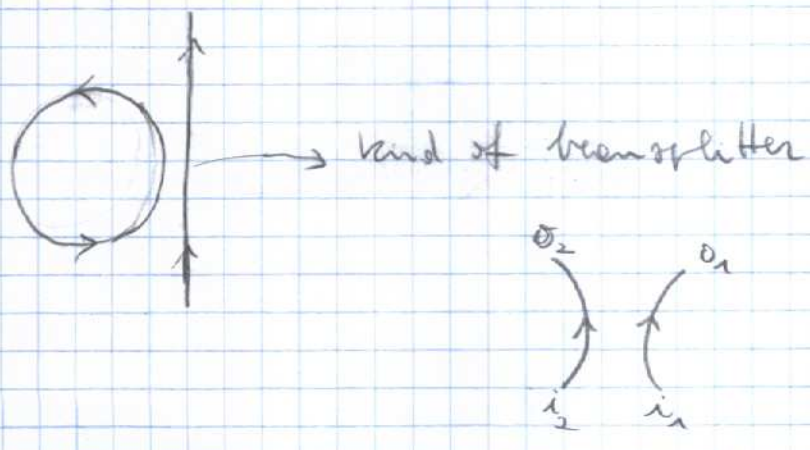


exercise 1 : fibers and rings

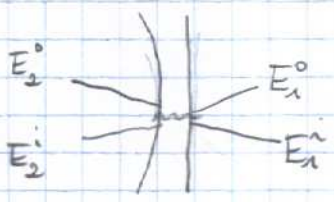


$$\begin{pmatrix} E_{o1} \\ E_{o2} \end{pmatrix} = M \begin{pmatrix} E_{i1} \\ E_{i2} \end{pmatrix}$$

Energy conservator $\rightarrow M^\dagger M = \mathbb{1}$

e.g. $M = \begin{pmatrix} t & r \\ -r & t \end{pmatrix}$, $|t|^2 + |r|^2 = 1$, $t, r \in \mathbb{R}$

Fiber Transmission



$$\left. \begin{aligned} E_2^o &= t E_2^i + r E_1^i \\ E_1^o &= t E_1^i + r E_2^i \end{aligned} \right\} \text{beam-splitter.}$$

$E_2^i = \exp(i k \cdot 2\pi R) E_2^o$: round-trip condition

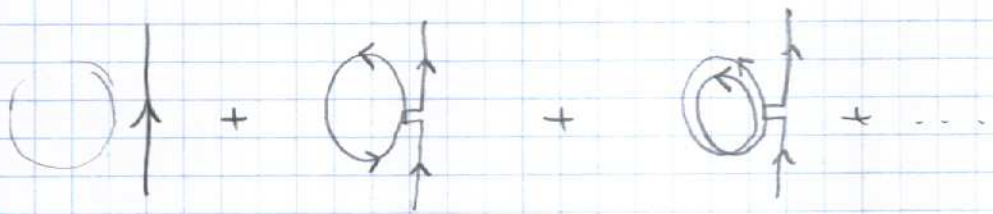
$$E_2^o = t \cdot \exp(i k 2lR) E_2^i = 2 E_1^i$$

$$E_2^o = \frac{-r}{1 - t \exp(i k 2lR)} E_1^i$$

$$E_1^o = t E_1^i - r \exp(i k 2lR) \frac{r}{1 - t \exp(i k 2lR)} E_1^i =$$

$$= E_1^i \left[t - \frac{r^2 e^{2i k l R}}{1 - t e^{2i k l R}} \right]$$

$$\approx E_1^i t + E_1^i (-r) e^{2i k l R} r + E_1^i (-r) e^{2i k l R} t e^{2i k l R} r + \dots$$



Non-absorbing limit: $k \in \mathbb{R}$, $|\exp(2i k l R)| = 1$

* resonant point $2\pi k R = 2\pi m$

$$E_1^o = E_1^i \left[t - \frac{1-t^2}{1-t} \right] = E_1^i [t - (1+t)] = -E_1^i$$

$$|E_1^o|^2 = |E_1^i|^2 \rightarrow \text{perfect transmission}$$

* in general :

$$E_1^o = E_1^i \left[t - \frac{(1-t^2) e^{2\pi i k R}}{1 - t e^{2\pi i k R}} \right] =$$

$$= E_1^i \left[\frac{t - t^2 e^{2\pi i k R} - e^{2\pi i k R} + t^2 e^{2\pi i k R}}{1 - t e^{2\pi i k R}} \right]$$

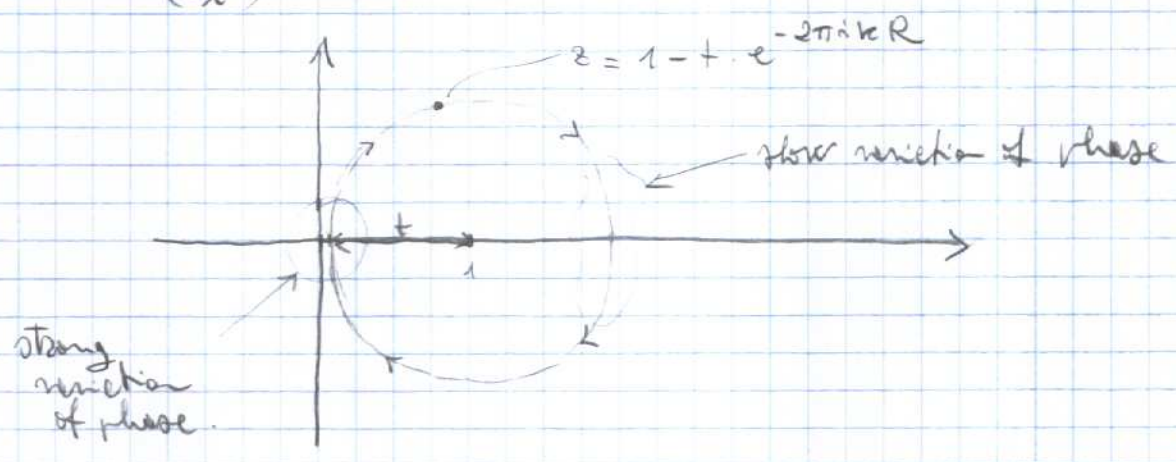
$$= -E_1^i e^{2\pi i k R} \left[\frac{1 - t e^{-2\pi i k R}}{1 - t e^{2\pi i k R}} \right]$$

as we have assumed $t \in \mathbb{R}$

$$\Rightarrow |E_1^o|^2 = |E_1^i|^2 \text{ for all frequencies!}$$

but transmission phase varies

$$\text{Arg} \left(\frac{E_1^o}{E_1^i} \right) = \pi + 2\pi k R + 2 \cdot \text{Arg} (1 - t e^{-2\pi i k R})$$



$$\frac{d}{dk} \text{Arg} \left(\frac{E_1^o}{E_1^i} \right) = \frac{2\pi R}{1-t} \approx \frac{8\pi R}{\epsilon^2}$$

Physical interpretation

$$\bar{\nu} = \frac{d}{d\omega} \phi(\omega) \approx \frac{1}{c} \frac{d}{dn} \phi(n) = \frac{8\pi R}{c} \cdot \frac{1}{2^2} =$$

$$= 4 \times T_{\text{round}} \times N_{\text{round}}$$

↘ # round trips before exiting
 (each round trip, probability 2^2 of exiting.)

Stationary-phase argument.

$$E_1^o(t) = \int d\omega E_1^i(\omega) \cdot t(\omega) e^{-i\omega t} \approx$$

$$\approx \int d\omega E_1^i(\omega) \cdot |t(\omega)| \cdot e^{i\phi(\omega)} e^{-i\omega t}$$

$$\approx \int d\omega E_1^i(\omega) |t(\omega)| e^{-i(\omega - \omega_0)(t - \frac{d\phi}{d\omega})} e^{-i\omega_0 t}$$

assuming $E_1^i(\omega)$ smooth and real

$$\Rightarrow E_1^o(t) \text{ peaked at } t - \frac{d\phi}{d\omega} = 0$$

$$\text{i.e. } t = \frac{d\phi}{d\omega}$$

on the other hand:

$$E_1^i(t) \approx \int d\omega E_1^i(\omega) e^{-i\omega t} \text{ peaked at } t=0.$$

$$\Rightarrow \text{DELAY TIME } \tau = \frac{d\phi}{d\omega}$$

In the presence of absorption:

$$t_{tot} = - \left[\frac{e^{2\pi i \tilde{n} R} - t}{1 - t e^{2\pi i \tilde{n} R}} \right]$$

where \tilde{n} = complex wavevector in ring

Single-round-trip loss $\exp[-2\pi R k_i]$

weak loss condition $2\pi R k_i \ll 1$

Non-radiative decay rate $\Gamma_{nr}/2 = \frac{c}{2\pi R} 2\pi R k_i = c k_i$

i.e. $\exp[-2\pi R k_i] = \exp[-2\pi R \Gamma_{nr}/2c]$

At resonance $\exp[2\pi i R n_r] = 1$

$$t_{tot} = \frac{t - e^{-2\pi R \Gamma_{nr}/2c}}{1 - t e^{-2\pi R \Gamma_{nr}/2c}} \approx \frac{t - 1 + 2\pi R \Gamma_{nr}/2c}{1 - t + t 2\pi R \Gamma_{nr}/2c}$$

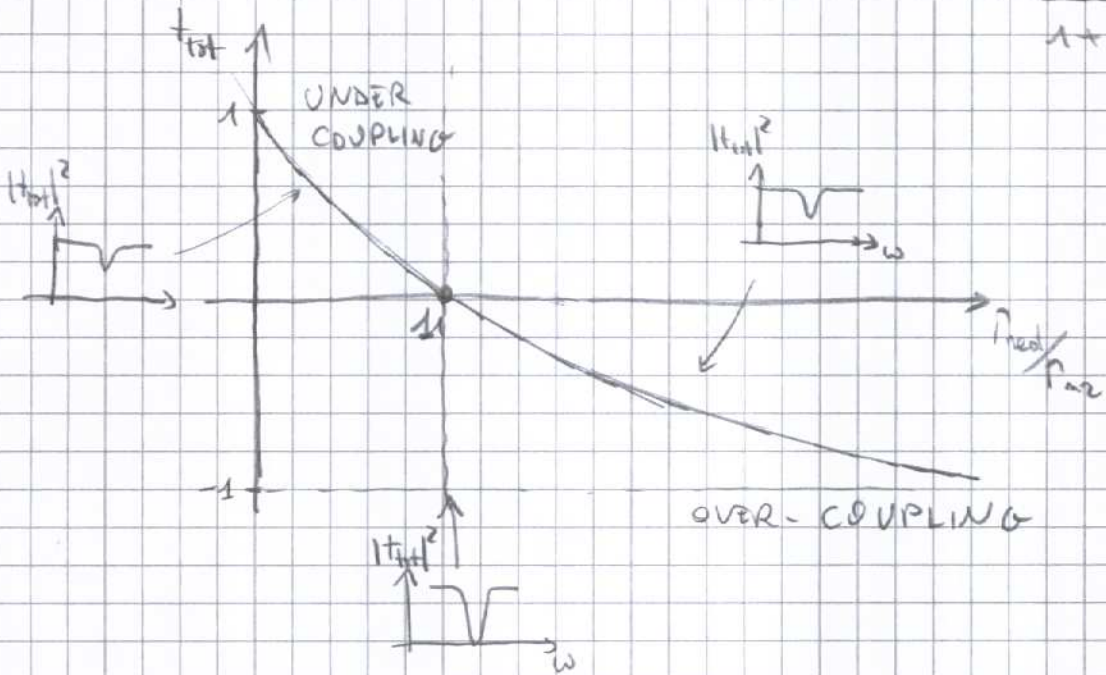
Radiative width:

denominator: $\frac{1}{1 - t e^{2\pi i \tilde{n} R}} \approx \frac{1}{1 - t - 2\pi i (n_r - n_i) R t}$

$$= \frac{1}{1 - t - 2\pi i R t \frac{\Delta \omega}{c}} = \frac{1}{i \frac{2\pi R t}{c} \Delta \omega - i \frac{c}{2\pi R t} (1 - t)}$$

$$\Rightarrow \frac{P_{rad}}{2} = \frac{c}{2\pi R t} (1 - t) \approx \frac{c}{2\pi R}$$

$$t_{tot} = \frac{-2\pi R/c P_{rad}/2 + 2\pi R/c P_{in}/2}{\frac{2\pi R}{c} \frac{P_{rad}}{2} + \frac{2\pi R}{c} \frac{P_{in}}{2}} = \frac{P_{in} - P_{rad}}{P_{rad} + P_{in}} = \frac{1 - P_{rad}/P_{in}}{1 + P_{rad}/P_{in}}$$



Amplitude injection of energy into my cavity

→ Peak depth maximum at OPTIMAL COUPLING $P_{in} = P_{rad}$ then decreases on both sides.

Addenda to solution of Ex. 1

Energy conservation $\Rightarrow \|S|v\rangle\|^2 = \||v\rangle\|^2 \quad \forall |v\rangle$

$\langle v | S^\dagger S | v \rangle = \langle v | v \rangle \quad \forall v$

$\Rightarrow S^\dagger S = \mathbb{1}$

Symmetry 1 \leftrightarrow 2:

$T^{-1} S T = S \quad \text{with} \quad T = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$

$T^{-1} S T = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} S_{11} & S_{12} \\ S_{21} & S_{22} \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \begin{pmatrix} S_{22} & S_{21} \\ S_{12} & S_{11} \end{pmatrix}$

$T^{-1} S T = S \Rightarrow \underbrace{S_{11}}_r = \underbrace{S_{22}}_r, \quad \underbrace{S_{12}}_t = \underbrace{S_{21}}_t$

Unitarity:

$S^\dagger S = \begin{pmatrix} r^* & t^* \\ t^* & r^* \end{pmatrix} \begin{pmatrix} r & t \\ t & r \end{pmatrix} = \mathbb{1}$

$\Rightarrow \begin{cases} |r|^2 + |t|^2 = 1 \\ 2r^*t + 2t^*r = 0 \end{cases}$

Write $t = |t| \cdot e^{i\theta} \Rightarrow 2r^* |t| e^{i\theta} + 2|t| e^{-i\theta} r = 0$

$\Rightarrow r = i |t| e^{i\theta} \quad \text{with} \quad |r|^2 = 1 - |t|^2$

Transmittivity of diode:

$$E_{in} = E_{01} e^{i2\pi R\omega/c}$$

$$E_{01} = \tilde{r}|r| e^{i\theta} E_{in} + |t| e^{i\theta} E_{r2} = e^{-i2\pi R\omega/c} E_{in}$$

$$E_{in} = \frac{|t| e^{i\theta} E_{r2}}{e^{-i2\pi R\omega/c} - \tilde{r}|r| e^{i\theta}}$$

$$E_{02} = \tilde{r}|r| e^{i\theta} E_{r2} + \frac{|t|^2 e^{2i\theta}}{e^{-i2\pi R\omega/c} - \tilde{r}|r| e^{i\theta}} E_{r2}$$

$$= \left[\tilde{r}|r| e^{i\theta} + \frac{|t|^2 e^{2i\theta} e^{i2\pi R\omega/c}}{1 - \tilde{r}|r| e^{i\theta} e^{i2\pi R\omega/c}} \right] E_{r2}$$

$$\approx \left[\tilde{r}|r| e^{i\theta} + |t|^2 e^{2i\theta} e^{i2\pi R\omega/c} + |t|^2 \tilde{r}|r| e^{3i\theta} e^{i4\pi R\omega/c} + \dots \right] \times E_{r2}$$

$$= \left[\frac{\tilde{r}|r| e^{i\theta} + |t|^2 e^{2i\theta} e^{i2\pi R\omega/c} + |t|^2 \tilde{r}|r| e^{3i\theta} e^{i4\pi R\omega/c}}{1 - \tilde{r}|r| e^{i\theta} e^{i2\pi R\omega/c}} \right] E_{r2} =$$

$$= e^{2i\theta} e^{i2\pi R\omega/c} \left[\frac{1 + \tilde{r}|r| e^{-i\theta} e^{-i2\pi R\omega/c}}{1 - \tilde{r}|r| e^{i\theta} e^{i2\pi R\omega/c}} \right] E_{r2}$$

$\tau(\omega)$

$$|\tau(w)| = 1 \quad \text{es} \quad \tau(w) = e^{i\phi} \quad \left(\frac{1+A}{1+A^*} \right)$$

$$\tau(w) = e^{i\phi(w)}$$

$$\phi(w) = 2\theta + 2\pi R w/c + \text{Arg} \left[\frac{1 + i|z| \dots}{1 - i|z| \dots} \right]$$

$$= 2\theta + 2\pi R w/c + 2 \text{Arg} \left[1 + i|z| e^{-i\theta} e^{-i2\pi R w/c} \right]$$

$$= 2\theta + 2\pi R w/c + 2 \text{Arg} \left[1 + |z| e^{i \left[\frac{\pi}{2} - \theta - 2\pi R w/c \right]} \right]$$

$$\frac{d\phi}{dw} = \frac{2\pi R}{c} + 2 \frac{d}{dw} \text{Arg} \left[1 + |z| \cos \lambda(w) + i|z| \sin \lambda(w) \right]$$

$$= \frac{2\pi R}{c} + 2 \frac{d}{dw} \arctan \left[\frac{|z| \sin \lambda(w)}{1 + |z| \cos \lambda(w)} \right] =$$

$$= \frac{2\pi R}{c} + 2 \frac{|z| [|z| + \cos \lambda]}{1 + |z|^2 + 2|z| \cos \lambda} \cdot \frac{d\lambda}{dw} =$$

$$= \frac{2\pi R}{c} + 2 \frac{|z| [|z| + \cos \lambda]}{1 + |z|^2 + 2|z| \cos \lambda} \cdot (-2\pi R/c) =$$

$$= \frac{2\pi R}{c} \cdot \left[1 - \frac{2|z|^2 + 2|z| \cos \lambda}{1 + |z|^2 + 2|z| \cos \lambda} \right] = \frac{2\pi R}{c} \left[\frac{1 + |z|^2 + 2|z| \cos \lambda - 2|z|^2 - 2|z| \cos \lambda}{1 + |z|^2 + 2|z| \cos \lambda} \right]$$

$$= \frac{2\pi R}{c} \frac{1 - |z|^2}{1 + |z|^2 + 2|z| \cos \lambda}$$

Max $\frac{d\phi}{d\omega}$ for $\cos\lambda = -1$

$$\Rightarrow \frac{d\phi}{d\omega}|_{\max} = \frac{2\pi R}{c} \frac{1 - |r|^2}{1 + |r|^2 - 2|r|}$$

$$= \frac{2\pi R}{c} \left[\frac{1 + |r|}{1 - |r|} \right] \approx \frac{4\pi R}{c[1 + |r|]}$$

Note: for $|r|=1 \Rightarrow \frac{d\phi}{d\omega} = 0 \rightarrow$ ring decoupled.

Definition of ring mode: $\Gamma = \frac{1}{\frac{2\pi R}{c}} \cdot (1 - |r|^2) \approx \frac{1 - |r|}{\frac{\pi R}{c}}$

$$\left(\tau_{\text{del}}|_{\max} = 4\Gamma^{-1} \right)$$

NOTE: delay is max for $\cos\lambda = -1$, i.e.

$$\frac{\pi}{2} - \theta - \frac{2\pi R\omega}{c} = (2M+1)\pi$$

$$\frac{2\pi R\omega}{c} = \frac{\pi}{2} + (2M+1)\pi = 2M\pi - \frac{\pi}{2} - \theta$$

reflective phase.

With dissipation:

$$e^{2\pi i R \omega / c} \rightarrow e^{2\pi i R \omega / c} e^{-2\pi R \Gamma_{me} / 2c}$$

Γ_{me} = non-radiative decay rate = $2c \cdot k_i$

$$\tau(\omega) = e^{i\omega t} \frac{e^{2\pi i R \omega / c} e^{-2\pi R \Gamma_{me} / 2c} e^{i\omega t} + i|z|}{1 - i|z| e^{i\omega t} e^{2\pi i R \omega / c} e^{-2\pi R \Gamma_{me} / 2c}}$$

$\rightarrow \in \mathbb{R}^+$ on resonance.

$$\tau / \omega_0 = i e^{i\omega t} \frac{|z| - i e^{i\omega t} e^{2\pi i R \omega / c} e^{-2\pi R \Gamma_{me} / 2c}}{1 - i e^{i\omega t} e^{2\pi i R \omega / c} e^{-2\pi R \Gamma_{me} / 2c} |z|}$$

$$\rightarrow \frac{|z| - e^{-2\pi R \Gamma_{me} / 2c}}{1 - |z| e^{-2\pi R \Gamma_{me} / 2c}} \approx$$

$$\approx \frac{|z| - 1 + \frac{2\pi R \Gamma_{me}}{2c}}{1 - |z| + \frac{2\pi R \Gamma_{me}}{2c} |z|} =$$

$$= \frac{-\frac{\pi R}{c} \Gamma_{rad} + \frac{\pi R}{c} \Gamma_{nr}}{\frac{\pi R}{c} \Gamma_{rad} + \Gamma_{nr} \frac{\pi R}{c}} = \frac{\Gamma_{nr} - \Gamma_{rad}}{\Gamma_{rad} + \Gamma_{nr}}$$